

Block Sliding Down a Frictionless Wedge

(Taylor Example 7.5) Consider a block with mass m sliding frictionlessly down an wedge with mass M that makes an angle α . The wedge itself slides frictionlessly across a horizontal floor near the surface of Earth. The block is released from the top of the wedge, with both objects initially at rest.

If length of the sloping face is d , how long does the block take to reach the bottom?

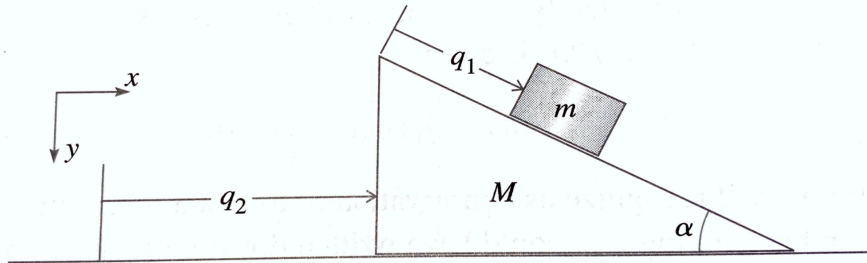


Figure 7.8 A block of mass m slides down a wedge of mass M , which is free to slide over the horizontal table.

Solution Since the generalized coordinates are defined in the picture, I'll start by writing down the Lagrangian:

$$\mathcal{L} = KE - PE$$

$$KE = \frac{1}{2}mv_m^2 + \frac{1}{2}Mv_M^2$$

Finding v_m is tricky, but here are two ways to do it:

Method 1: Use the Cartesian basis vectors.

$$\begin{aligned} \vec{r}_m &= (q_2 + q_1 \cos \alpha)\hat{x} + q_1 \sin \alpha \hat{y} \\ \vec{v}_m &= \frac{d}{dt}\vec{r} = (\dot{q}_2 + \dot{q}_1 \cos \alpha)\hat{x} + \dot{q}_1 \sin \alpha \hat{y} \\ v_m^2 &= \vec{v} \cdot \vec{v} \\ &= (\dot{q}_2 + \dot{q}_1 \cos \alpha)^2 + (\dot{q}_1 \sin \alpha)^2 \\ &= \dot{q}_2^2 + \dot{q}_1^2 \cos^2 \alpha + 2\dot{q}_1\dot{q}_2 \cos \alpha + \dot{q}_1^2 \sin^2 \alpha \\ &= \dot{q}_2^2 + \dot{q}_1^2 (\cos^2 \alpha + \sin^2 \alpha) + 2\dot{q}_1\dot{q}_2 \cos \alpha \\ v_m^2 &= \dot{q}_2^2 + \dot{q}_1^2 + 2\dot{q}_1\dot{q}_2 \cos \alpha \end{aligned}$$

Method 2: Use non-orthogonal basis vectors.

$$\vec{r}_m = q_2 \hat{q}_2 + q_1 \hat{q}_1$$

where \hat{q}_2 and \hat{q}_1 are normal but not orthogonal:

$$\begin{aligned}\hat{q}_2 \cdot \hat{q}_2 &= 1 \\ \hat{q}_1 \cdot \hat{q}_1 &= 1 \\ \hat{q}_2 \cdot \hat{q}_1 &= \cos \alpha\end{aligned}$$

Now take the derivative to find the velocity:

$$\begin{aligned}\vec{v}_m &= \frac{d}{dt} \vec{r} = \dot{q}_2 \hat{q}_2 + \dot{q}_1 \hat{q}_1 \\ v_m^2 &= \vec{v} \cdot \vec{v} \\ &= (\dot{q}_2 \hat{q}_2 + \dot{q}_1 \hat{q}_1) \cdot (\dot{q}_2 \hat{q}_2 + \dot{q}_1 \hat{q}_1)\end{aligned}$$

Multiplying (or FOILing) I get:

$$\begin{aligned}v_m^2 &= \dot{q}_2^2 (\hat{q}_2 \cdot \hat{q}_2) + 2 \dot{q}_2 \dot{q}_1 (\hat{q}_2 \cdot \hat{q}_1) + \dot{q}_1^2 (\hat{q}_1 \cdot \hat{q}_1) \\ v_m^2 &= \dot{q}_2^2 + \dot{q}_1^2 + 2 \dot{q}_1 \dot{q}_2 \cos \alpha\end{aligned}$$

Finding v_M is more straightforward because the coordinate q_2 completely describes its position:

$$v_M^2 = \dot{q}_2^2$$

So, the kinetic energy is:

$$KE = \frac{1}{2} m (\dot{q}_2^2 + \dot{q}_1^2 + 2 \dot{q}_1 \dot{q}_2 \cos \alpha) + \frac{1}{2} M \dot{q}_2^2$$

The potential energy is near Earth potential energy and only the box changes potential energy:

$$PE = -mgq_1 \sin \alpha$$

So, the Lagrangian function is:

$$\begin{aligned}\mathcal{L} &= KE - PE \\ \mathcal{L} &= \frac{1}{2} m (\dot{q}_2^2 + \dot{q}_1^2 + 2 \dot{q}_1 \dot{q}_2 \cos \alpha) + \frac{1}{2} M \dot{q}_2^2 + mgq_1 \sin \alpha\end{aligned}$$

Now, apply the Euler-Lagrange Equation for each coordinate:

$$\begin{aligned}\frac{\partial \mathcal{L}}{\partial q_1} &= \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{q}_1} \\ mg \sin \alpha &= \frac{d}{dt} (m\dot{q}_1 + m\dot{q}_2 \cos \alpha) \\ \therefore g \sin \alpha &= \ddot{q}_1 + \ddot{q}_2 \cos \alpha\end{aligned}$$

$$\begin{aligned}\frac{\partial \mathcal{L}}{\partial q_2} &= \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{q}_2} \\ 0 &= \frac{d}{dt} (m\dot{q}_2 + m\dot{q}_1 \cos \alpha + M\dot{q}_2) \\ 0 &= (m + M)\ddot{q}_2 + m\ddot{q}_1 \cos \alpha \\ \therefore \ddot{q}_2 &= \frac{m \cos \alpha}{m + M} \ddot{q}_1\end{aligned}$$

These are coupled 2nd order differential equations, but they are linear and I can decouple them by substitution:

$$\begin{aligned}g \sin \alpha &= \ddot{q}_1 + \left(\frac{m \cos \alpha}{m + M} \ddot{q}_1 \right) \cos \alpha \\ \ddot{q}_1 &= \frac{g \sin \alpha}{1 - \frac{m \cos^2 \alpha}{m + M}} \\ &= \frac{g \sin \alpha (m + M)}{m + M - m \cos^2 \alpha} \\ \ddot{q}_1 &= \frac{g \sin \alpha (m + M)}{M + m(1 - \cos^2 \alpha)} \checkmark \\ \ddot{q}_2 &= \left(\frac{m \cos \alpha}{m + M} \right) \left(\frac{g \sin \alpha (m + M)}{M + m(1 - \cos^2 \alpha)} \right) \\ \ddot{q}_2 &= \frac{mg \sin \alpha \cos \alpha}{M + m(1 - \cos^2 \alpha)} \checkmark\end{aligned}$$

Do some evaluative sensemaking (Does my answer make sense?)

- Dimensions - both accelerations have dimensions of acceleration (or Length/Time²) ✓

- Functional Behavior - if I think about the system as a whole, the net force on the system is due to the acceleration due to gravity and the normal force from the floor on the wedge. These forces are constant - therefore, I expect the accelerations to be constant - and they are!

Are there any momenta conserved? Yes! Since the net force on the system is vertical, I expect the horizontal momentum to be conserved. q_2 is missing from the Lagrangian, so the momentum $\partial\mathcal{L}/\partial\dot{q}_2$ is constant in time, i.e., conserved. This looks like:

$$M\dot{q}_2 + m(\dot{q}_2 + \dot{q}_1 \cos \alpha) = \text{constant}$$

- Special Cases - there are a bunch we could try! Here are a few.

1. Wedge is Flat: $\alpha \rightarrow 0$

Expectation: Both accelerations are zero - boring.

Check:

$$\ddot{q}_1 = \frac{g \sin(0)(m + M)}{M + m(1 - \cos^2(0))} = 0 \checkmark$$

$$\ddot{q}_2 = \frac{mg \sin(0) \cos(0)}{M + m(1 - \cos^2(0))} = 0 \checkmark$$

2. Wedge is vertical: $\alpha \rightarrow \pi/2$

Expectation: Box is in freefall, wedge doesn't move.

Check:

$$\ddot{q}_1 = \frac{g \sin(\pi/2)(m + M)}{M + m(1 - \cos^2(\pi/2))} = g \checkmark$$

$$\ddot{q}_2 = \frac{mg \sin(\pi/2) \cos(\pi/2)}{M + m(1 - \cos^2(\pi/2))} = 0 \checkmark$$

3. Wedge is much heavier than the box: $M \gg m$

Expectation: Box sliding down an inclined plane. Wedge doesn't move, box has constant acceleration $g \sin \alpha$.

Check:

$$\ddot{q}_1 = \frac{g \sin \alpha (m + M)}{M + m(1 - \cos^2 \alpha)} \stackrel{\rightarrow}{=} g \sin \alpha \checkmark$$

$$\ddot{q}_2 = \frac{mg \sin \alpha \cos \alpha}{M + m(1 - \cos^2 \alpha)} \stackrel{\rightarrow}{=} 0 \checkmark$$

4. Box is much heavier than wedge: $m \gg M$

Expectation: Wedge will shoot out from under the box really fast. Box will essentially fall straight down.

Check:

$$\ddot{q}_1 = \frac{g \sin \alpha (m + M)}{M + m(1 - \cos^2 \alpha)} = \frac{g}{\sin \alpha} ??$$

$$\ddot{q}_2 = \frac{mg \sin \alpha \cos \alpha}{M + m(1 - \cos^2 \alpha)} = \frac{g}{\tan \alpha} ??$$

This is a tough one to interpret. It really helps to convert the q 's into Cartesian coordinates:

$$\ddot{y} = \ddot{q}_1 \cos \alpha$$

$$= g$$

As for \ddot{q}_2 , this depend on alpha. As $\alpha \rightarrow \pi/2$, then the acceleration go to zero. This makes sense because the normal force from the box on the wedge is small the steeper the wedge is. As $\alpha \rightarrow 0$, the normal force is bigger, and I see that the acceleration get really big.

Now, find the time! (Don't forget the finish answering the question.)

Since the acceleration is constant, I can use my knowledge of kinematics for constant acceleration to find the time:

$$q_1(t) = q_1(0) + \dot{q}_1(0)t + \frac{1}{2}\ddot{q}_1 t^2$$

$$\rightarrow t = \sqrt{\frac{2q_1}{\ddot{q}_1}}$$

$$t = \sqrt{\frac{2d[M + m(1 - \cos^2 \alpha)]}{g \sin \alpha (m + M)}}$$

Some sensemaking about the time:

- Length of incline: d increases

Expectation: The time will increase.

Check: $t \propto \sqrt{d}$ ✓

- Compare to fixed incline:

Expectation: Time will be shorter for movable wedge: $\frac{t_{movable}}{t_{fixed}} < 1$.

Check: For fixed wedge, $a = g \sin \alpha \rightarrow t_{fixed} = \sqrt{\frac{2d}{g \sin \alpha}}$

$$\begin{aligned} \frac{t_{movable}}{t_{fixed}} &= \sqrt{\frac{\frac{2d[M+m(1-\cos^2 \alpha)]}{g \sin \alpha (m+M)}}{\frac{2d}{g \sin \alpha}}} \\ &= \sqrt{\frac{M+m(1-\cos^2 \alpha)}{m+M}} \\ &\leq 1 \checkmark \text{ because } (1-\cos^2 \alpha) \leq 1 \end{aligned}$$